

Plateau's Problem

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1 Motivation: Minimal Surfaces in Nature

When a metal wire is dipped in a soap solution, **surface tension** forces the resulting film to be of minimal surface area; a soap film spanning a wire loop minimizes surface area subject to that boundary.

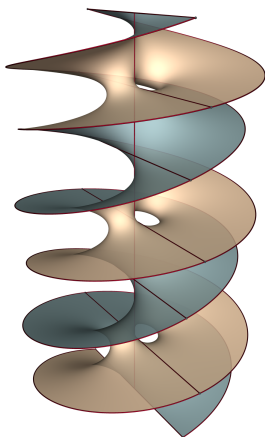
Plateau's Problem, first raised by Joseph-Louis Lagrange in 1760, claims that there exists such a minimal surface for every simple wire loop; soap films are an example of this.

We can now extend this to a mathematical context. The metal wire becomes our **Jordan curve** Γ , while our soap film becomes the desired minimal surface.

Note that all parts of the problem covered in this paper will be solely within \mathbb{R}^3 .

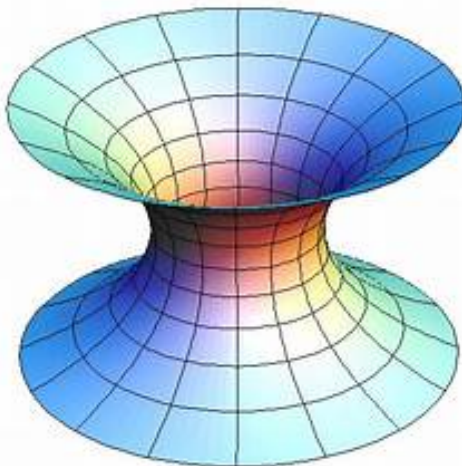
Some examples of minimal surfaces within \mathbb{R}^3 :

- **Helicoid**: The minimal surface formed along a helix:



A helicoid is parametrized as $\sigma(u, v) = (v \cos u, v \sin u, u)$.

- **Catenoid:** The minimal surface formed along a catenary:



A catenoid is parametrized as $\sigma(u, v) = (\cosh v \cos u, \cosh v \sin u, v)$.

There are many other examples of minimal surfaces, but these are the most common.

2 Plateau's Problem: Statement

For some Jordan curve $\Gamma \in \mathbb{R}^3$, some surface S exists, such that:

1. S **spans** the Jordan curve: $\partial S = \Gamma$,
2. S **minimizes** its area among all surfaces spanning Γ : $A(S) \leq A(S')$ for any other surface S' .

Note: A (classical) minimal surface is one with mean curvature $H \equiv 0$ in the interior. Area-minimizing implies minimality, but not conversely.

For the sake of simplicity, we will not discuss self-intersecting curves.

3 Plateau's Problem: Douglas's Solution

Let D be the **closed unit disk**:

$$D = \{(x, y) \in \mathbb{R}^2 | x^2 + y^2 \leq 1\}.$$

Fix a **continuous, one-to-one, counterclockwise** parametrization:

$$\gamma : S^1 \rightarrow \Gamma.$$

Our goal is to find a surface of minimal area that spans Γ .

When we try to minimize directly over maps $u : D \rightarrow \mathbb{R}^3$ with $u|_{\partial D} = \gamma \circ \phi$, the boundary reparametrization $\phi : S^1 \rightarrow S^1$ creates a **huge symmetry**: composing with disk automorphisms leads to a lack of compactness.

Douglas's fix: restricting to **orientation-preserving homeomorphisms** ϕ that fix 3 boundary points (e.g. $1, e^{2i\pi/3}, e^{4i\pi/3}$). This 3-point normalization helps us further work towards the existence of a minimizer.

Given ϕ , we define our boundary map:

$$f_\phi := \gamma \circ \phi : S^1 \rightarrow \mathbb{R}^3.$$

Let $u_\phi : D \rightarrow \mathbb{R}^3$ be the unique **harmonic extension** (Poisson integral) of f_ϕ .

Note that among all maps with the same boundary data f_ϕ , u_ϕ minimizes the **Dirichlet energy**:

$$E[u] := \frac{1}{2} \int_D |\nabla u|^2 dx dy = \frac{1}{2} \int_D (|u_x|^2 + |u_y|^2) dx dy.$$

Now, note that given **conformal** u , area is equivalent to energy:

$$A[u] = E[u],$$

as:

$$A[u] = \int_D |\partial_x u \times \partial_y u| = \frac{1}{2} \int_D |\nabla u|^2 \text{ when } u \text{ is conformal.}$$

Thus, if we can confirm that our energy-minimizing harmonic extension u is indeed conformal, then we will have our area minimizer.

The Douglas Functional: Douglas discovered a boundary-only functional $D(\phi) = E[u_\phi]$, which measures a boundary roughness of f_ϕ . Thanks to our equality, we note that minimizing $D(\phi)$ over ϕ is equivalent to minimizing the interior energy among harmonic fillings of reparametrized boundaries.

For reference, a form of Douglas's functional is the integral:

$$D(\phi) = \frac{1}{4\pi} \int_0^{2\pi} \int_0^{2\pi} \frac{|\gamma(\phi(e^{it})) - \gamma(\phi(e^{is}))|^2}{2(1 - \cos(t - s))} dt ds.$$

The derivation of the formula is too complex for this paper, but a short derivation using Fourier modes of the Poisson extension shows $E[u_\phi] = D(\phi)$.

Douglas also showed that $D(\phi)$ is **bounded below** and **lower semicontinuous** on the normalized class of ϕ 's, as well as the fact that any minimizing sequence ϕ_k has a **convergent subsequence**.

Thus, there exists a minimizer ϕ_* , such that:

$$D(\phi_*) = \inf_{\phi} D(\phi).$$

Now, let $u := u_{\phi_*}$ be the harmonic map with boundary f_{ϕ_*} .

Observe the first variation of D under smooth one-parameter perturbations ϕ_t , with $\phi_0 = \phi_*$. Stationarity of D at ϕ_* translates, after pulling through the Poisson extension, into **boundary conditions** for the harmonic map u :

$$\langle u_r, u_\theta \rangle|_{\partial D} = 0 \text{ and } |u_r| = |u_\theta| \text{ on } \partial D,$$

where r, θ are the polar coordinates on our disk. These conditions state that the boundary principal directions are **orthogonal**, and that the boundary metric coefficients are **equal**, respectively.

Thus, our energy minimizer u is both harmonic and conformal.

Finally, we must relate conformality to **area minimality**. For any competitor $v : D \rightarrow \mathbb{R}^3$ with the same boundary trace as u :

$$A[v] \geq E[v] \geq E[u] = A[u],$$

where:

- $A[v] \geq E[v]$, and $A[v] = E[v] \iff v$ is conformal,
- $E[v] \geq E[u]$ because u is the harmonic energy minimizer for the boundary data,
- $A[u] = E[u]$ because u is conformal.

Therefore, u has the least area among all spanning maps with the given boundary trace. This proves existence of a solution to Plateau's problem for Γ \square

References

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